

Sensors, Signals and Noise

COURSE OUTLINE

- Introduction
- Signals and Noise: 3) Analysis and Simulation
- Filtering
- Sensors and associated electronics



Noise Analysis and Simulation

- White Noise
- Band-Limited White Noise (Wide-Band Noise)
- Basic Parameters of Wide-Band Noise
- Foundations of White-Noise Filtering
- Generation and Simulation of Any Noise by a Poisson Process
and for those who wish to gain a better insight
- APPENDIX: Noise Power Transients



White Noise



White Noise (stationary)

IDEAL «white» noise

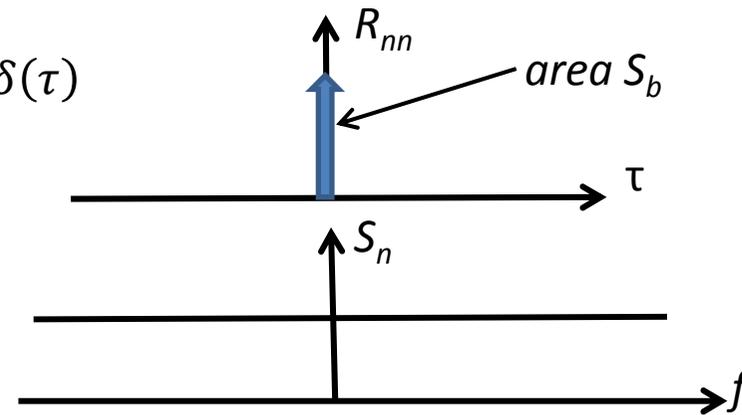
is a concept extrapolated from Johnson noise and shot noise
defined by its **essential** feature:

no autocorrelation at any time distance τ , no matter how small

$$R_{nn}(\tau) = S_b \cdot \delta(\tau)$$

and therefore constant spectrum

$$S_n(f) = S_b$$



In reality such a noise does not exist: it would have divergent power $\overline{n^2} \rightarrow \infty$

REAL «white» noise has

- **Very small width** of autocorrelation, shorter than the minimum time interval of interest in the actual case and therefore approximated to zero
- **Very wide band** with constant spectral density S_b , wider than the maximum frequency of interest in the actual case and therefore approximated to infinite



White Noise (non-stationary)

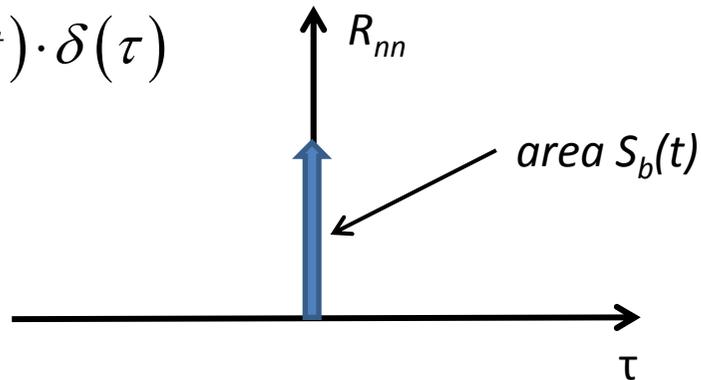
Also in non-stationary cases the IDEAL «white» noise is defined by the essential characteristic feature:

no correlation at any finite time distance τ , no matter how small, but the noise intensity is no more constant, it varies with time t

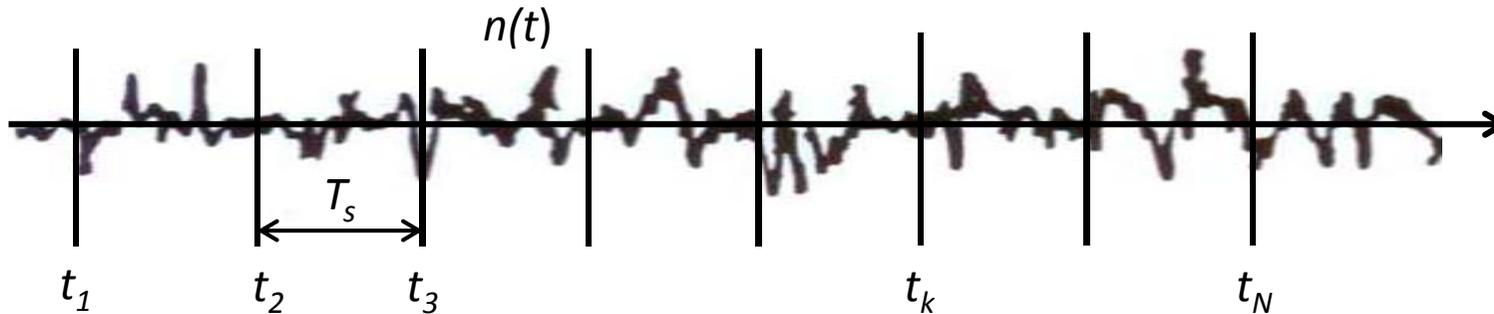
that is

the autocorrelation function is δ -like, but has time-dependent area $S_b(t)$

$$R_{nn}(t, t + \tau) = S_b(t) \cdot \delta(\tau)$$



Filtering white noise is simple



For clarity, let's consider a discrete case:

linear filtering in digital signal processing:

- Sample n_1 at t_1 and multiply by a weight w_1 ,
- sample n_2 at $t_2 = t_1 + T_s$ and multiply by a weight w_2 and sum
- and so on

The filtered noise n_f is

$$n_f = w_1 n_1 + w_2 n_2 + \dots = \sum_{k=1}^N w_k n_k$$

and its mean square value is

$$\begin{aligned} \overline{n_f^2} &= \overline{w_1^2 n_1^2 + w_2^2 n_2^2 + \dots + w_1 n_1 \cdot w_2 n_2 + w_1 n_1 \cdot w_3 n_3 + \dots} = \\ &= w_1^2 \overline{n_1^2} + w_2^2 \overline{n_2^2} + \dots + w_1 w_2 \overline{n_1 n_2} + w_1 w_3 \overline{n_1 n_3} + \dots \end{aligned}$$



Filtering white noise is simple

$$\begin{aligned}\overline{n_f^2} &= \overline{w_1^2 n_1^2 + w_2^2 n_2^2 + \dots + w_1 n_1 \cdot w_2 n_2 + w_1 n_1 \cdot w_3 n_3 + \dots} = \\ &= \overline{w_1^2 n_1^2} + \overline{w_2^2 n_2^2} + \dots + \overline{w_1 w_2 n_1 n_2} + \overline{w_1 w_3 n_1 n_3} + \dots\end{aligned}$$

If noise at interval T_s is **not correlated**, then all rectangular terms vanish

$$\overline{n_1 n_2} = \overline{n_1 n_2} = \dots = 0$$

and the result is **simply a sum of squares**, even in case of non-stationary noise

$$\begin{aligned}\overline{n_f^2} &= \overline{w_1^2 n_1^2} + \overline{w_2^2 n_2^2} + \dots = \\ &= \sum_{k=1}^N \overline{w_k^2 n_k^2}\end{aligned}$$

If the noise is stationary

$$\overline{n_1^2} = \overline{n_2^2} = \overline{n_3^2} = \dots = \overline{n^2}$$

there is a further simplification

$$\begin{aligned}\overline{n_f^2} &= \overline{n^2} (w_1^2 + w_2^2 + \dots) = \\ &= \overline{n^2} \sum_{k=1}^N w_k^2\end{aligned}$$

we will see later that also with continuous filtering white noise brings similar simplification



Band-Limited White Noise or Wide-Band Noise

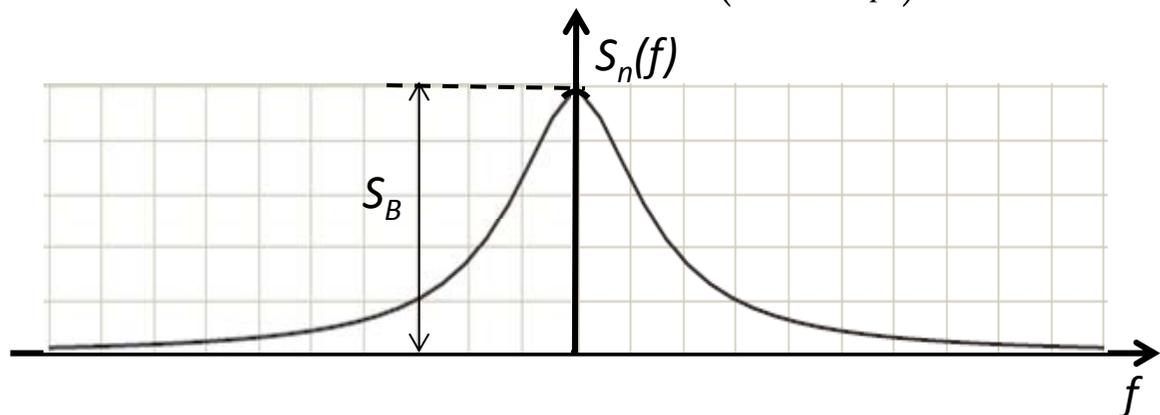
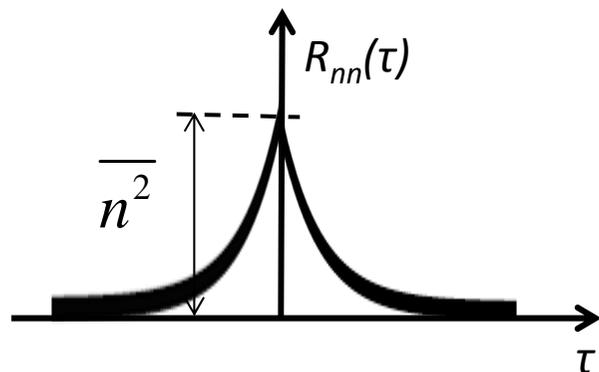


Band-limited white noise (wide-band noise)

- Real white noise = white noise with band limited at high frequency
- The limit may be inherent in the noise source or due to low-pass filtering enforced by the circuitry. Anyway, in all real cases there is such a limit
- A frequent typical case is the **Lorentzian** spectrum:
band limited by a **simple pole** with time constant T_p , pole frequency $f_p = 1/2\pi T_p$

$$R_{nn}(\tau) = \overline{n^2} e^{-\frac{|\tau|}{T_p}}$$

$$S_n(f) = \frac{S_B}{1 + (2\pi f T_p)^2}$$



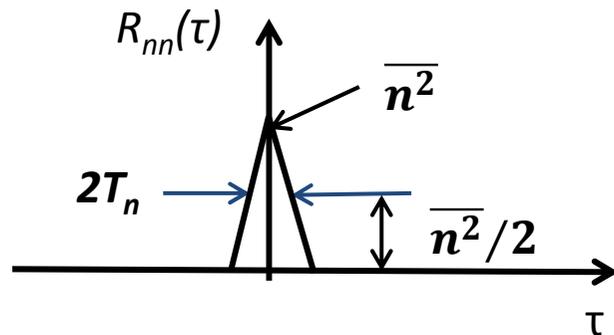
Basic Parameters of Wide-Band Noise



Simplified description of wide-band noise

The true $R_{nn}(\tau)$ and $S_n(f)$ can be **approximated** by simple functions retaining the noise main features:

- a) equal mean square $\overline{n^2}$ and b) equal spectral density S_b



in time: $R_{nn}(\tau)$ triangular approx, half-width $2T_n$

a) equal msq noise : $R_{nn}(0) = \overline{n^2}$

b) equal spectral density: [area of $R_{nn}(\tau)$] = S_b ,
i.e. $\overline{n^2} 2 T_n = S_b$

Correlation width = area/peak

$$\Delta\tau = 2T_n$$

in frequency: $S_n(f)$ rectang approx, half-width f_n

a) equal msq noise : [area of $S_n(f)$] = $\overline{n^2}$
i.e. $S_b 2 f_n = \overline{n^2}$

b) equal spectral density: $S_n(0) = S_b$

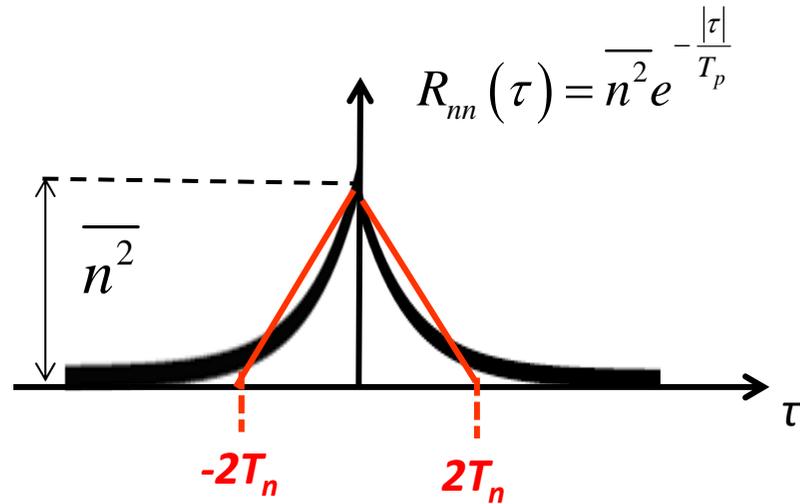
Noise bandwidth: area/peak

$$\Delta f = 2f_n$$

Note that $\Delta\tau \cdot \Delta f = 1$ which is consistent with $S_n(f) = F[R_{nn}(\tau)]$

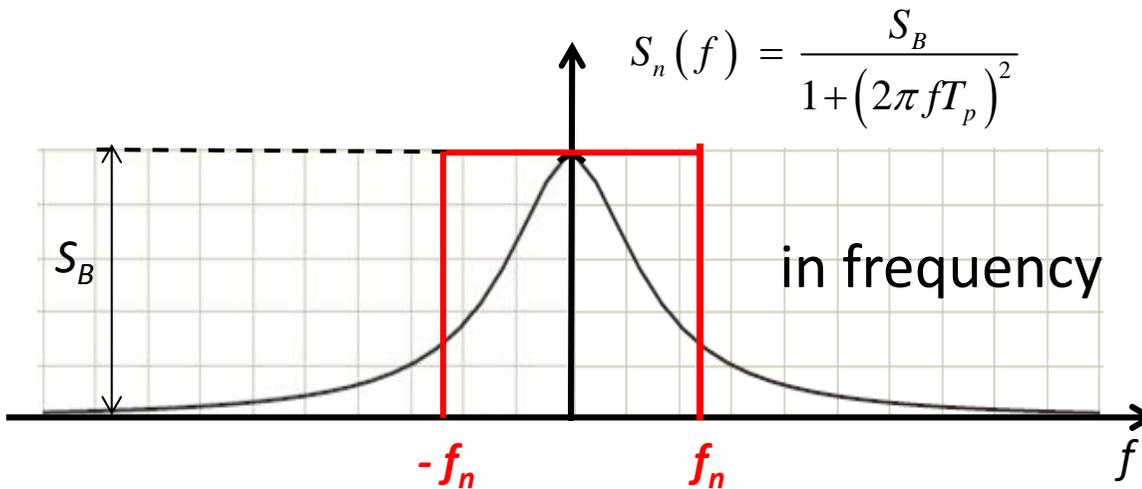


Simplified description of Lorentzian spectrum



$$S_B = \int_{-\infty}^{\infty} R_{nn}(\tau) d\tau = \overline{n^2} 2T_p$$

$$T_n = T_p$$



$$\overline{n^2} = \int_{-\infty}^{\infty} S_n(f) df = S_B \frac{1}{2T_p}$$

$$2f_n = \frac{1}{2T_p}$$

Note that $f_n \neq f_p$, namely $f_n = \frac{1}{4T_p} = \frac{\pi}{2} f_p$



Foundations of White-Noise Filtering

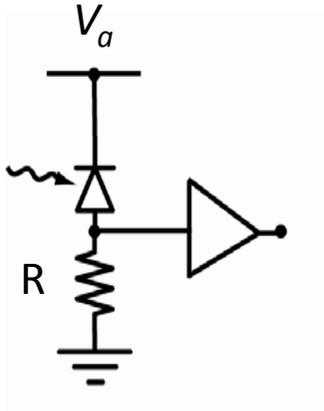


Noise filtering clarified by the Poisson pulse model

- Noise is a random superposition of elementary pulses
- The elementary pulse type (i.e. pulse waveform and its F-transform) defines the noise type (i.e. autocorrelation shape and spectrum shape)
- The passage through a linear constant-parameter filter modifies the elementary pulse type
- The pulse modification causes a corresponding modification of the noise
- Noise filtering can thus be understood, studied and evaluated by understanding, studying and evaluating the filtering of the elementary pulses



Low-pass filtering of White noise



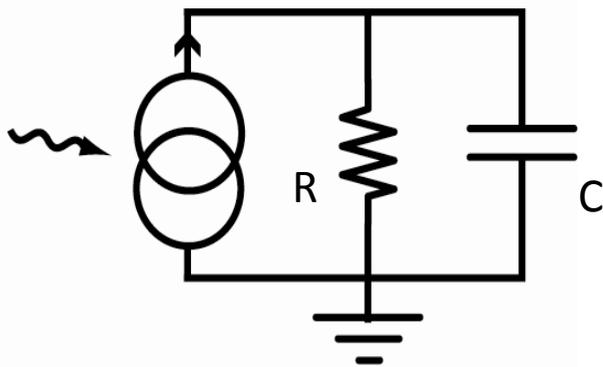
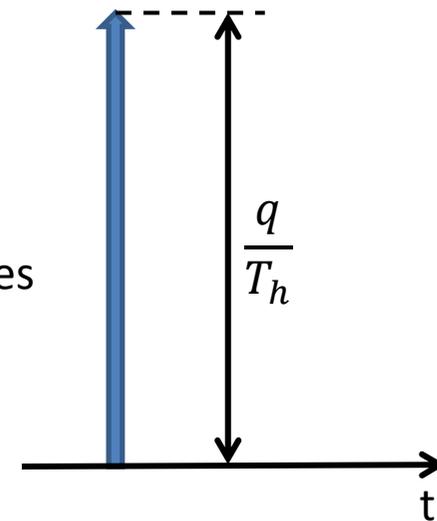
shot current white noise:

$$S_i(f) \cong qI \quad \text{for } f \ll 10\text{GHz}$$

$$R_{ii}(\tau) \cong qI\delta(\tau) \quad \text{for } \tau \gg 100\text{ps}$$

diode current: elementary short pulses
with rate $p = I/q$, $T_h \cong 100\text{ps}$
approximate δ -pulses

$$qh(t) \cong q\delta(t)$$



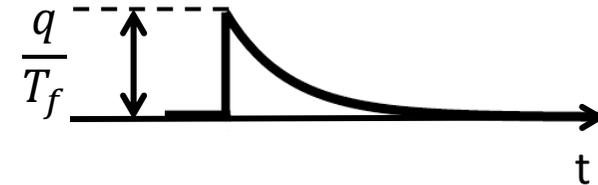
current in R: elementary exponential pulses

with rate $p = I/q$, $T_f = RC$

$$q \cdot f(t)$$

$$f(t) = \frac{1}{T_f} 1(t)e^{-t/T_f}$$

$$F(f) = \frac{1}{1 + j2\pi f T_f}$$



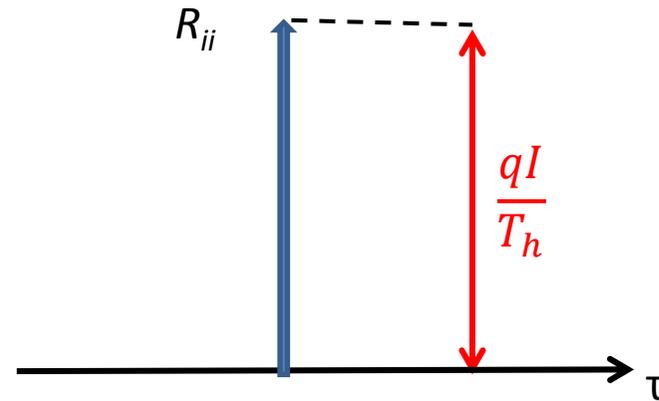
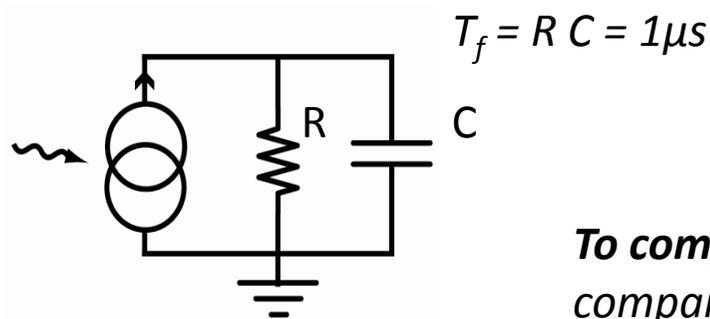
e.g. with $R = 100\text{k}\Omega$ and $C = 10\text{pF}$ we have $T_f = 1\mu\text{s}$



Low-pass filtering of White noise: time domain view

Input noise (current in the diode):
 δ -like autocorrelation (width $\approx 100ps$)

$$R_{ii}(\tau) \cong qI\delta(\tau) \quad \text{for } \tau \gg 100ps$$

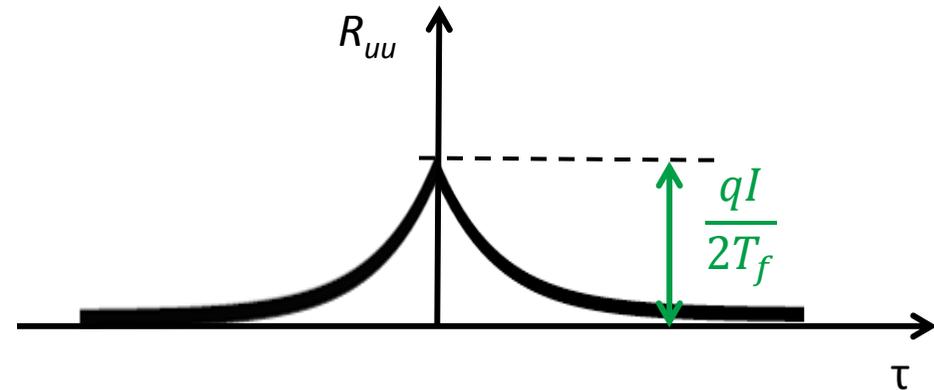


*To compare msq values of noise **before** and **after** filtering compare the **central values** of autocorrelation functions*

Output noise (current in R):
 autocorrelation function

$$R_{uu}(\tau) = qI \cdot k_{ff}(\tau)$$

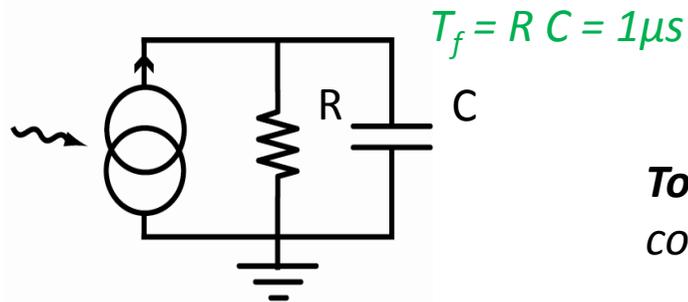
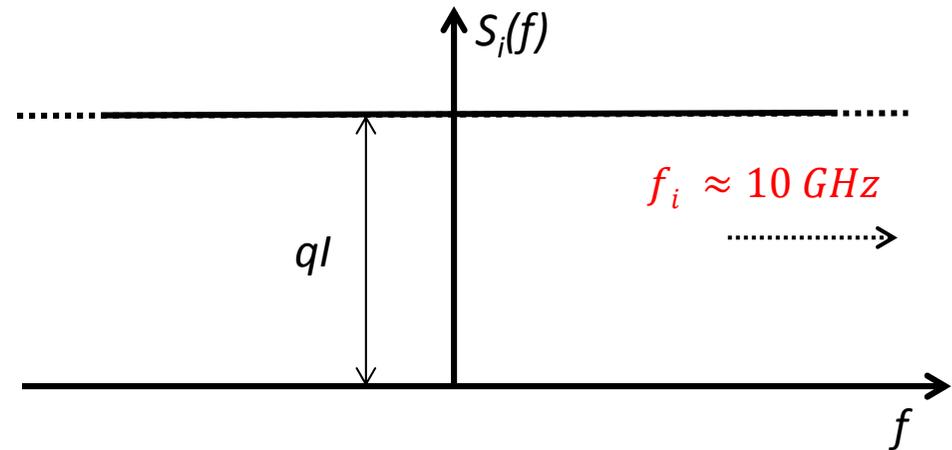
$$k_{ff}(\tau) = \frac{1}{2T_f} e^{-\left|\frac{\tau}{T_f}\right|}$$



Low-pass filtering of White noise: frequency domain view

Input noise (diode current):
spectral density S_i constant
(bandwidth $f_i \approx 10 \text{ GHz}$)

$$S_i(f) = S_b \quad \text{for } f \ll 10 \text{ GHz}$$

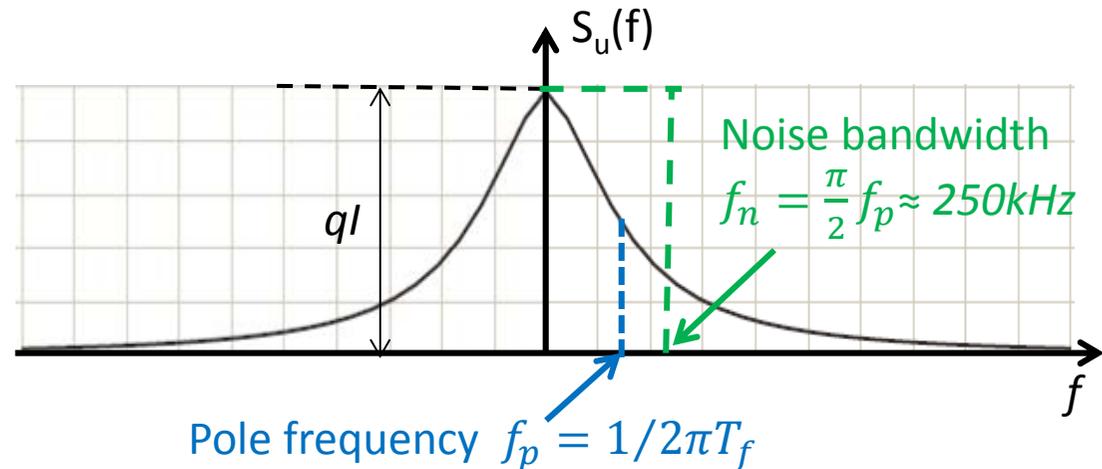


To compare msq values of noise **before** and **after** filtering
compare the **areas** of input and output spectral densities

Output noise (current in R):
spectral density function $S_u(f)$

$$S_u(f) = qI \cdot |F(f)|^2$$

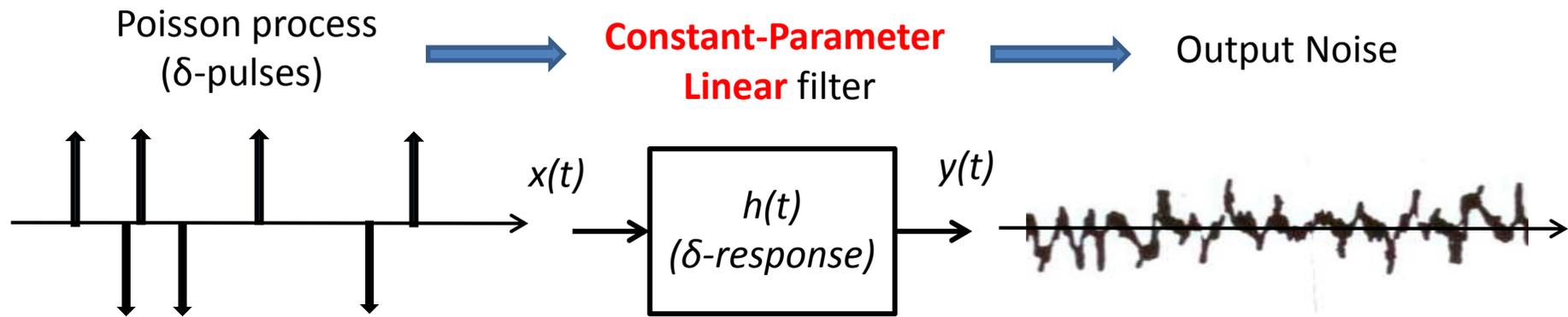
$$|F(f)|^2 = \frac{1}{1 + (2\pi f T_f)^2}$$



Generation and Simulation of Any Noise by a Poisson Process



Generation of any noise from Poisson process



p pulse rate (p/s)
 Q pulse area

Filter autocorrelation

$$k_{hh}(\tau) = \int_{-\infty}^{\infty} h(t)h(t + \tau)dt$$

Transfer Function

$$H(f) = F[h(t)]$$

White Input Noise

$$\begin{cases} R_{xx}(\tau) = pQ^2 \delta(\tau) \\ S_x(f) = pQ^2 \end{cases}$$

Output Noise

$$\begin{cases} R_{yy}(\tau) = pQ^2 k_{hh}(\tau) \\ S_y(f) = pQ^2 |H(f)|^2 \end{cases}$$



Generation of any noise from Poisson process

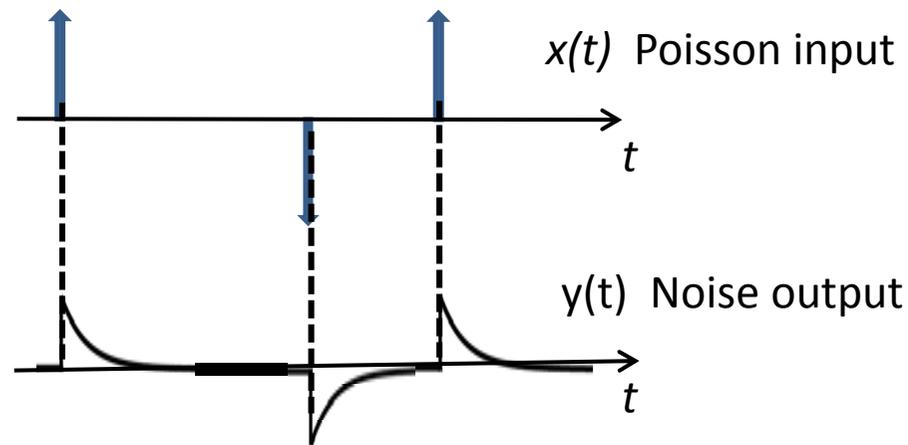
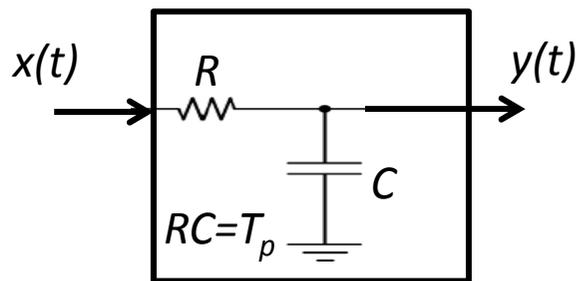
For producing a given $S_y(f)$ (i.e. a given $R_{yy}(\tau) = F^{-1}[S_y(f)]$) the filter must have

- $|H(f)|^2 = S_y(f)$ normalized to 1 at $f=0$
- $k_{hh}(\tau) = R_{yy}(\tau)$ normalized to unit area

Example: band-limited white noise $S_y = S_B \frac{1}{1 + (f/f_p)^2}$ with $f_p = \frac{1}{2\pi T_p}$

$$H(\omega) = \frac{1}{1 + j\omega T_p}$$

$$h(t) = \frac{1}{T_p} e^{-\frac{t}{T_p}} 1(t)$$



$$S_y(f) = pQ^2 |H(f)|^2 = pQ^2 \frac{1}{1 + (f/f_p)^2}$$



Appendix:

Noise Power Transients

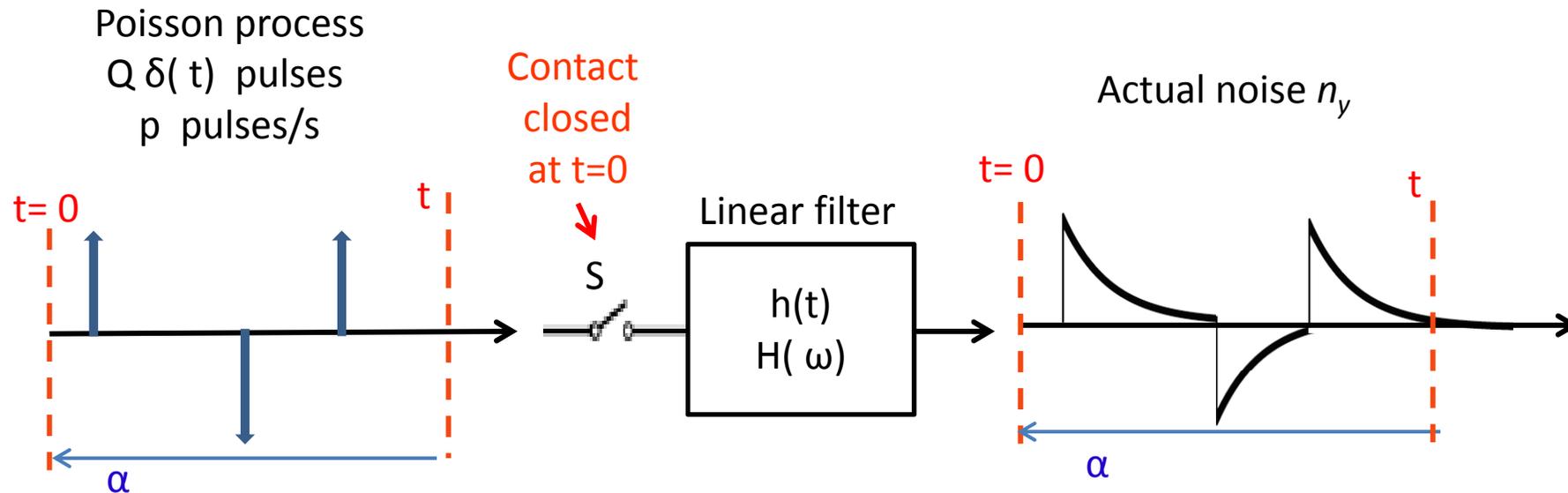
Q: how does the noise power rise when a noise source is switched on ?

A: the Poisson model clarifies!



Rise-time of the Noise power

Noise modeling by Poisson process also shows how noise rises after switch-on



Noise switch-on at $t=0$ modeled by closing at $t=0$ the switch at the filter input.

The mean square output at time t is computed by Campbell's theorem

but integrating **only over the interval where output pulses occur**, i.e. for $0 < \alpha < t$,

$$\overline{n_y^2(t)} = pQ^2 \int_0^t h^2(\alpha) d\alpha$$



Rise-time of real white noise

The rise of the output noise intensity can be observed by computing the root-mean square noise (rms) versus time t (normalized to unit spectral density of the input Poisson noise)

$$\sigma_y(t) = \sqrt{\frac{n_y^2(t)}{pQ^2}} = \sqrt{\int_0^t h^2(\alpha) d\alpha}$$

Let's consider real White Noise with band-limit due to a simple pole with time constant T_p

- The noise is modeled by pulses $h(t) = \frac{1}{T_p} e^{-\frac{t}{T_p}} 1(t)$
- The rise of the rms noise is

$$\sigma_y(t) = \sqrt{\int_0^t \frac{1}{T_p^2} e^{-\frac{2\alpha}{T_p}} d\alpha} = \sqrt{\frac{1}{2T_p} \cdot \left(1 - e^{-\frac{2t}{T_p}}\right)}$$



Rise-time of real white noise

$$\sigma_y(t) = \sqrt{\int_0^t \frac{1}{T_p^2} e^{-\frac{2\alpha}{T_p}} d\alpha} = \sqrt{\frac{1}{2T_p} \cdot \left(1 - e^{-\frac{2t}{T_p}}\right)}$$

Wide-band «white» Noise

- The time constant T_p is short and the intensity (rms) rises swiftly, reaching in a few T_p the steady value

$$\sigma_y = \sqrt{\frac{1}{2T_p}} \quad \text{for } t \gg T_p$$

Moderately Wide-band Noise

- the time constant T_p is longer and the intensity rise is slower.

In the first part $t \ll T_p$ and the intensity rises approximately as \sqrt{t}

$$\sigma_y(t) = \sqrt{\frac{1}{2T_p} \cdot \left(1 - e^{-\frac{2t}{T_p}}\right)} \approx \frac{\sqrt{t}}{T_p}$$



Rise-time of random-walk noise

«Random-Walk» Noise»

- denotes noise with spectral density $\propto 1/\omega^2$.
- is generated by integration of white noise
(e.g. shot current noise $S_i = 2qI$ in capacitor $C \rightarrow$ voltage noise $S_v = 2qI / C^2 \omega^2$)
- is modeled by step elementary pulses $h(t) = 1(t)$
- The variance rises as \sqrt{t}

$$\sigma_y(t) = \sqrt{\int_0^t h^2(\alpha) d\alpha} = \sqrt{\int_0^t d\alpha} = \sqrt{t}$$

- At long time $t \rightarrow \infty$ the variance is divergent
In the frequency domain this corresponds to the power over a band extended down to very low frequency $f \rightarrow 0$

$$\sigma_y^2 = \lim_{f_i \rightarrow 0} \int_{f_i}^{\infty} |H(f)|^2 df = \lim_{f_i \rightarrow 0} \int_{f_i}^{\infty} \frac{1}{f^2} df \rightarrow \infty$$



Rise-time of random-walk noise

